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SPECTROSCOPICAL EFFECTS ARISING UNDER APPLICATION
OF PULSE SUPPLY TO ZINC HOLLOW CATHODE DISCHARGE

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ABSTRACT

The intensity and width of the Zn 330,2 nm spectral line emitted from hollow cathode discharge in pulse mode depending on pulse parameter changes are investigated. The purpose of this study is to obtain the effect of pulse mode supply and to find the optimum values of pulse parameters for stable emission of intensive as well as narrow zinc spectral lines. The line intensity gain is maximum 100, while the line width is less or equal to that obtained at the d.c.mode.

INTRODUCTION

Application of pulse supply to the hollow cathode discharge increases its advantages: stable emission of intensive as well as narrow spectral lines^{1,2,3,4}. It is considered that the effect from the application of pulse supply is predominantly due to decreased discharge sensitivity towards the cathode thermal regime. This implies that the pulse supply will be considerably useful for the investigation of elements like zinc, easily entering the discharge zone. Zinc is characterized

by a high sputtering rate⁵, low evaporation temperature and high vapor pressure⁶. As a result its atoms enter very intensively in the discharge zone at d.c. mode which leads to discharge instabilities and broadening of the emitted spectral lines⁷. Investigations on zinc discharge in pulse mode indicate that the Zn 201,3 nm spectral line intensity can be enhanced⁸ and the quality of zinc atomic absorption analysis-improved⁹.

In the present paper the intensity I and width $\Delta\lambda$ of the Zn 330,2 nm spectral line excited in a hollow cathode discharge in pulse mode as a function of pulse parameter changes (amplitude A , width t and pulse period T) are investigated. The results are compared with those ($I_0, \Delta\lambda_0$) obtained from discharge in d.c. mode for the same value of the mean current. The purpose of this study is to estimate the advantages of pulse mode supply and to find the optimum values of pulse parameters for stable emission of intensive as well as narrow zinc spectral lines emitted from a hollow cathode discharge in pulse mode.

EXPERIMENTAL

Fig. 1 shows our experimental set-up. The discharge tube is supplied with monopolar rectangular pulses and d.c. component not exceeding 2 mA. Pulse parameters are changed to its maximum values at which the discharge is still stable and arc mode does not appear. The hollow cathode discharge tube consists of a replaced cylindrical cathode and ring shaped anode put at the distance 1 mm from the cathode front. The cathode length is 15 mm and its inner diameter is 3 mm. Pure zinc as cathode material is not suitable because of the very intensive entering of zinc atoms in the discharge zone, which

leads to the discharge instability at lower values of pulse parameters and easier appearing of arc mode. That's why the alloy 50% Zn + 50% Al was used. The pressure of neon in the discharge tube is about 1,5 torr. The profile of the investigated Zn 330,2 nm spectral line was recorded using a scanning Fabry-Perot interferometer by changing the air pressure in the barocamera.

RESULTS AND DISCUSSION

Figs.2 and 3 represent the behaviour of the function $I/I_0 = f(\alpha, t, T)$. The following conclusions can be made having in mind the figures: 1) I/I_0 depends at the same time on the three pulse parameters (α, t, T) ; 2) $I/I_0 = f(t)$ is most pronounced; 3) $I/I_0 = f(\alpha)$ is more strongly expressed at lower values of T ; 4) The maximum value of I/I_0 - about 100 - almost does not change at the increasing of T from 220 μ s to 1597 μ s but only shifts

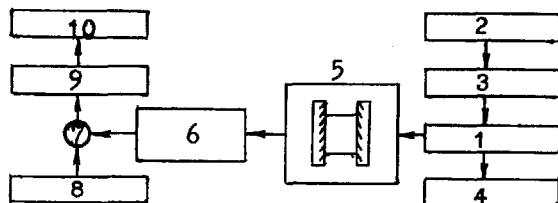


FIG.1

Experimental set-up principal scheme: 1 - discharge tube with a hollow cathode; 2 - pulse generator; 3 - pulse amplifier; 4 - oscillograph; 5 - Fabry-Perot interferometer in the barocamera; 6 - mono-chromator DMR-4; 7 - photomultiplier FEU (39,106); 8 - photomultiplier supply; 9 - d.c. amplifier; 10 - recorder G1B1.

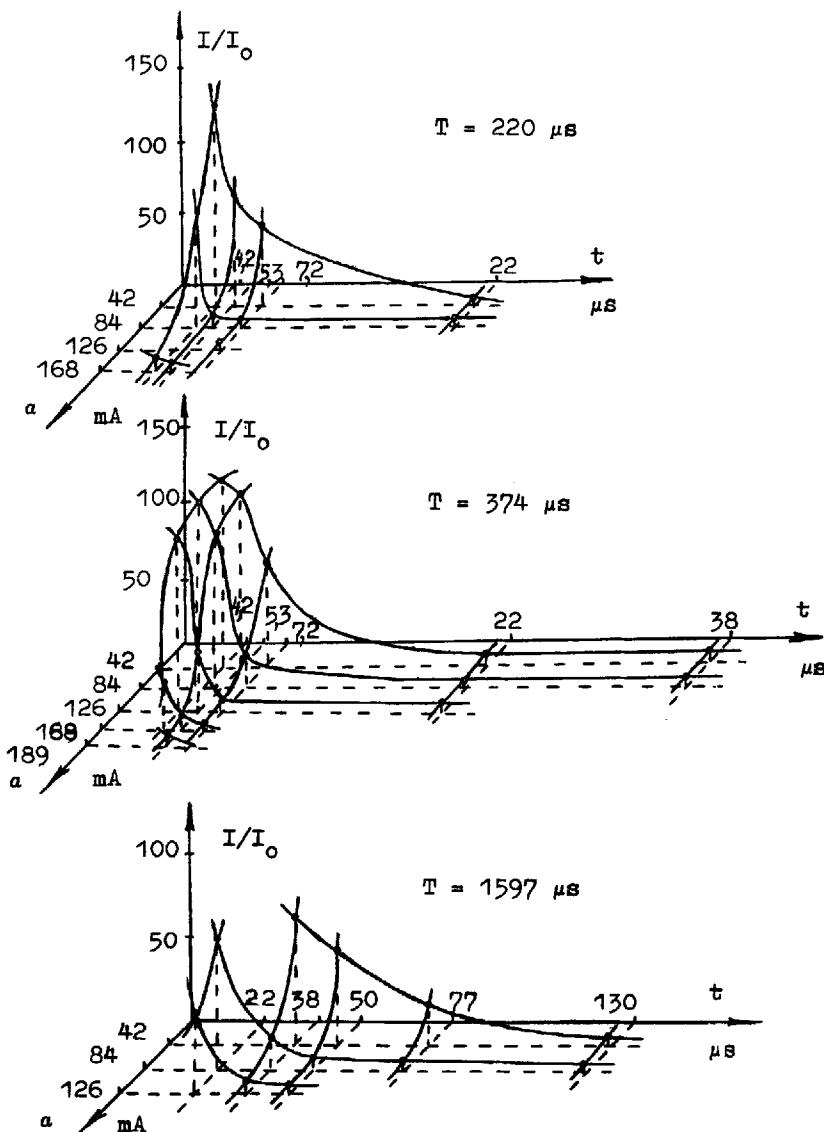


FIG.2

Dependence of Zn 330,2 nm spectral line intensity increase I/I_0 on the parameters of pulse supply.

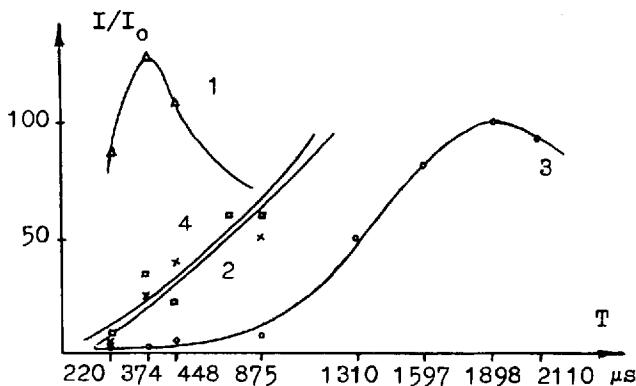


FIG.3

Dependence of Zn 330,2 nm spectral line intensity increase I/I_0 on T : 1- $t=4,2 \mu s$ at $a=84$ mA; 2- $t=7,2 \mu s$ at $a=84$ mA; 3- $t=22 \mu s$ at $a=84$ mA; 4- $t=4,2 \mu s$ at $a=168$ mA.

to higher values of t and a ; 5) I/I_0 decreases at $T \geq 1597 \mu s$. In tables 1 and 2 the widths of the Zn 330,2 nm spectral line emitted from dc. and pulse mode discharge correspondingly are summarized.

The Gaussian ($\Delta\nu_G$, $\Delta\nu_{G0}$) and the Dispersion ($\Delta\nu_D$, $\Delta\nu_{D0}$) components of the recorded widths ($\Delta\nu$, $\Delta\nu_0$) are obtained according to Ballik method¹⁰. Corrections for Gaussian type broadening (due to the round diaphragm $\delta\nu_1$ and to the surface mirror defects $\delta\nu_2$) and of Dispersion type (as a result of the mirror reflection $R - \delta\nu_R$) were made by the equations given in:¹¹

$$\delta\nu_1 = \frac{\lambda r^2}{\pi F^2}, \quad \delta\nu_2 = \frac{4l}{\pi}, \quad \delta\nu_R = \frac{1-R}{2nlF^2},$$

where r is the radius of the round diaphragm, F is the focus distance of the collimating lens, l is the distance between the mirrors. Having in mind the corrections in our case ($\delta\nu_1 = 15$ mK, $\delta\nu_2 \approx 24$ mK, $\delta\nu_R = 24$ mK at $R = 90\%$, $l = 0,7$ sm, $r = 0,05$ sm, $F = 50$ sm, $\Delta l \approx \lambda/60$)

we obtain the proper Gaussian ($\delta\nu_G$, $\delta\nu_{Go}$) and the proper Dispersion ($\delta\nu_D$, $\delta\nu_{Do}$) components. Tables 1 and 2 indicate that:
 1) $\delta\nu_D$ and $\delta\nu_{Do}$ do not depend on pulse parameters a , t and T ;
 2) $\delta\nu_G$ and $\delta\nu_{Go}$ are 5-10 times higher than $\delta\nu_D$ and $\delta\nu_{Do}$. This coincides with the results obtained for the predominantly Gaussian character of the profile of the lines emitted from hollow cathode discharge¹¹. Besides it can be concluded that application of pulse supply gives rise mainly to thermal changes in the discharge. More clearly the dependence $\delta\nu_G = f(t, T)$ at $a = 168$ mA and $\delta\nu_{Go} = f(i_o)$ /dashed line/ is shown on fig. 4. Having in mind the results in tables and fig. 4, the following conclusions can be made: 1) $\delta\nu_G$ increases with a and t increasing, this increase is faster at $a > 100$ mA; 2) $\delta\nu_G$ at $T = 220$ μ s, $T = 374$ μ s and $T = 448$ μ s has approximately equal values, smaller or equal

TABLE 1

Dependence of Zn 330,2 nm widths $\Delta\nu_o$, $\Delta\nu_{Go}$, $\Delta\nu_{Do}$, $\delta\nu_{Go}$, $\delta\nu_{Do}$ [mK] on the discharge current i_o .

i_o [mA]	$\Delta\nu_o$	$\Delta\nu_{Go}$	$\Delta\nu_{Do}$	$\delta\nu_{Go}$	$\delta\nu_{Do}$
3,1	186,0	140,0	46,0	136,5	22,0
3,7	204,0	151,5	52,0	148,5	28,0
4,1	210,0	164,0	46,0	161,5	22,0
5,5	217,5	173,0	46,0	170,7	22,0
6,6	226,5	181,0	46,0	178,8	22,0
7,6	238,0	190,0	49,3	200,0	25,3
8,6	248,5	202,0	49,3	200,0	25,3
10,0	258,5	214,0	46,0	212,0	22,0

TABLE 2

Dependence of Zn 330,2 nm spectral line widths $\Delta\delta$ _G, $\Delta\delta$ _D, $\delta\delta$ _G, $\delta\delta$ _D [mK] on the pulse supply parameters (α , t , T).

T μs	α mA	4,2				5,3				7,2				22				
		$\Delta\delta$ _G	$\Delta\delta$ _D															
220	42	155511304451092205	163012094451162205	172313404601310220	197516094901577160													
	84	172013204001290160	175013504451320205	17861370451342215	197016204151596175													
	126	175013404451310205	180514304151403175	183214504281427188														
	168	190514804601457220	195515204151523175	220017404601723220														
374	42	150310704451030205	161211504601114220	174313204451290205	195515204451497205													
	84	160011904151157175	173813104451280205	177413504451320205	169916094901574160													
	126	168812704151238175	179513904281360188	182214304281403188														
	168	189514704451443205	194015104451483205	215017004601672220														
448	168	189014604601434220	193015004601473220	2180117104761687216														
	875	190014604451434205	193014904451460205	209016704281643188	24652030476201216													
	168																	
	2299	168																

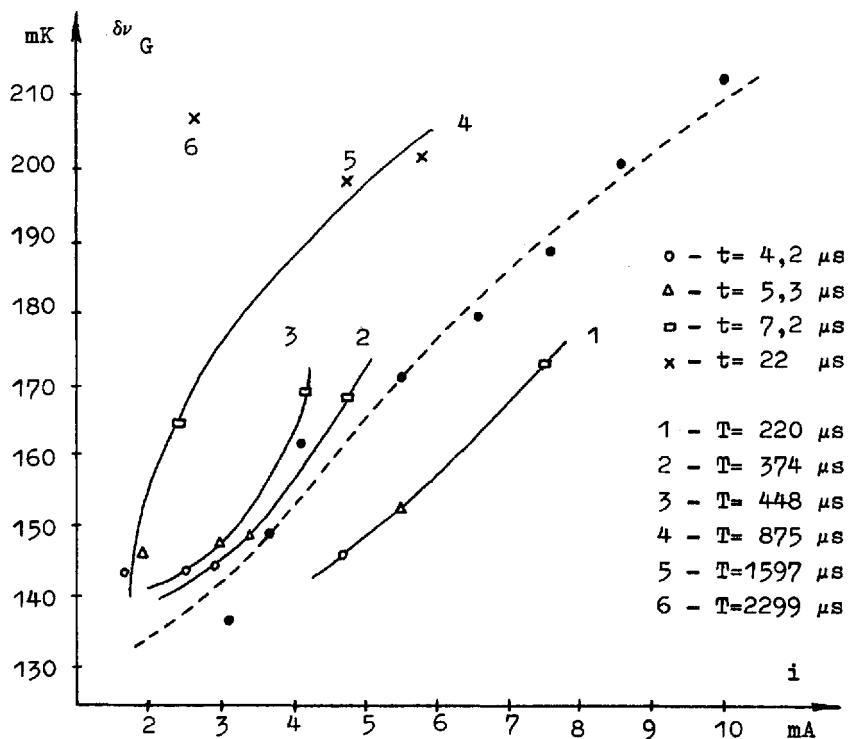


FIG.4

Dependences of proper Gaussian components $\delta y_G = f(t, T)$ at $a = 168$ mA and $\delta y_{G_0} = (i_0)$ /dashed line/ at the corresponding mean current values

to δy_{G_0} ; 3) δy_G at $T \geq 875$ μ s increases and becomes bigger than δy_{G_0} . While the first result can be explained by Doppler broadening and self-absorption because of mean current increasing in the discharge, such explanation for the second and the third results is not satisfactory. On the other hand our observations on the discharge at $T \geq 875$ μ s show that the shining diffused volume between the cathode and anode begins

to increase. Drawing a general conclusion from all these experimental results we can assume that at the given conditions the effect of the directed movement of the atoms from the cathode cavity outwards is manifested. Atomic transport phenomena from the cathode have been observed by other authors too^{12,13,14}. Detailed study of the investigated zinc line profile can lead to a more precise conclusion concerning the possibility of the assumed transport phenomena. The qualitative differences between our results and those for other elements obtained by the above cited authors regardless of some distinctions in the discharge tube constructions is predominantly due, in our opinion to the different properties of the investigated elements.

Generalizing the results obtained for the intensity and width of the Zn 330,2 nm spectral line emitted from hollow cathode discharge for our discharge tube construction in pulse mode, we can conclude that the optimum conditions for stable emission of zinc lines with increased intensity and at the same time with widths smaller or equal to that obtained at the d.c. mode are: $t=4,2+5,3 \mu s$, $a=40+80$ mA and $T=220+448 \mu s$.

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