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SPECTROSCOPICAL EFFECTS ARISING UNDER APPLICATION
OF PULSE SUPPLY TO ZINC HOLLOW CATHODE DISCHARGE

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ABSTRACT

The intensity and width of the Zn 330,2 nm spectral line emitted from hollow cathode discharge in pulse mode depending on pulse parameter changes are investigated. The purpose of this study is to obtain the effect of pulse mode supply and to find the optimum values of pulse parameters for stable emission of intensive as well as narrow zinc spectral lines. The line intensity gain is maximum 100, while the line width is less or equal to that obtained at the d.c. mode.

INTRODUCTION

Application of pulse supply to the hollow cathode discharge increases its advantages: stable emission of intensive as well as narrow spectral lines^{1,2,3,4}. It is considered that the effect from the application of pulse supply is predominantly due to decreased discharge sensitivity towards the cathode thermal regime. This implies that the pulse supply will be considerably useful for the investigation of elements like zinc, easily entering the discharge zone. Zinc is characterized

by a high sputtering rate⁵, low evaporation temperature and high vapor pressure⁶. As a result its atoms enter very intensively in the discharge zone at d.c. mode which leads to discharge instabilities and broadening of the emitted spectral lines⁷. Investigations on zinc discharge in pulse mode indicate that the Zn 201,3 nm spectral line intensity can be enhanced⁸ and the quality of zinc atomic absorption analysis-improved⁹.

In the present paper the intensity I and width $\Delta\lambda$ of the Zn 330,2 nm spectral line excited in a hollow cathode discharge in pulse mode as a function of pulse parameter changes (amplitude U , width t and pulse period T) are investigated. The results are compared with those $(I_0, \Delta\lambda_0)$ obtained from discharge in d.c. mode for the same value of the mean current. The purpose of this study is to estimate the advantages of pulse mode supply and to find the optimum values of pulse parameters for stable emission of intensive as well as narrow zinc spectral lines emitted from a hollow cathode discharge in pulse mode.

EXPERIMENTAL

Fig. 1 shows our experimental set-up. The discharge tube is supplied with monopolar rectangular pulses and d.c. component not exceeding 2 mA. Pulse parameters are changed to its maximum values at which the discharge is still stable and arc mode does not appear. The hollow cathode discharge tube consists of a replaced cylindrical cathode and ring shaped anode put at the distance 1 mm from the cathode front. The cathode length is 15 mm and its inner diameter is 3 mm. Pure zinc as cathode material is not suitable because of the very intensive entering of zinc atoms in the discharge zone, which

leads to the discharge instability at lower values of pulse parameters and easier appearing of arc mode. That's why the alloy 50% Zn + 50% Al was used. The pressure of neon in the discharge tube is about 1,5 torr. The profile of the investigated Zn 330,2 nm spectral line was recorded using a scanning Fabry-Perot interferometer by changing the air pressure in the barocamera.

RESULTS AND DISCUSSION

Figs. 2 and 3 represent the behaviour of the function $I/I_0 = f(\alpha, t, T)$. The following conclusions can be made having in mind the figures: 1) I/I_0 depends at the same time on the three pulse parameters (α, t, T); 2) $I/I_0 = f(t)$ is most pronounced; 3) $I/I_0 = f(\alpha)$ is more strongly expressed at lower values of T ; 4) The maximum value of I/I_0 - about 100 - almost does not change at the increasing of T from 220 μs to 1597 μs but only shifts

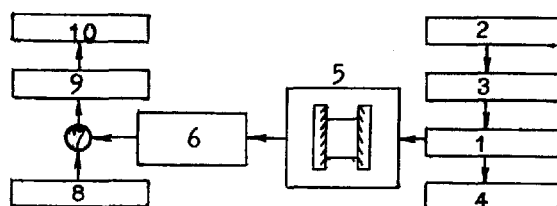


FIG. 1

Experimental set-up principal scheme: 1-discharge tube with a hollow cathode; 2-pulse generator; 3-pulse amplifier; 4-oscilloscope; 5-Fabry-Perot interferometer in the barocamera; 6-monochromator DMR-4; 7-photomultiplier FEU (39,106); 8-photomultiplier supply; 9-d.c. amplifier; 10-recorder G1B1.

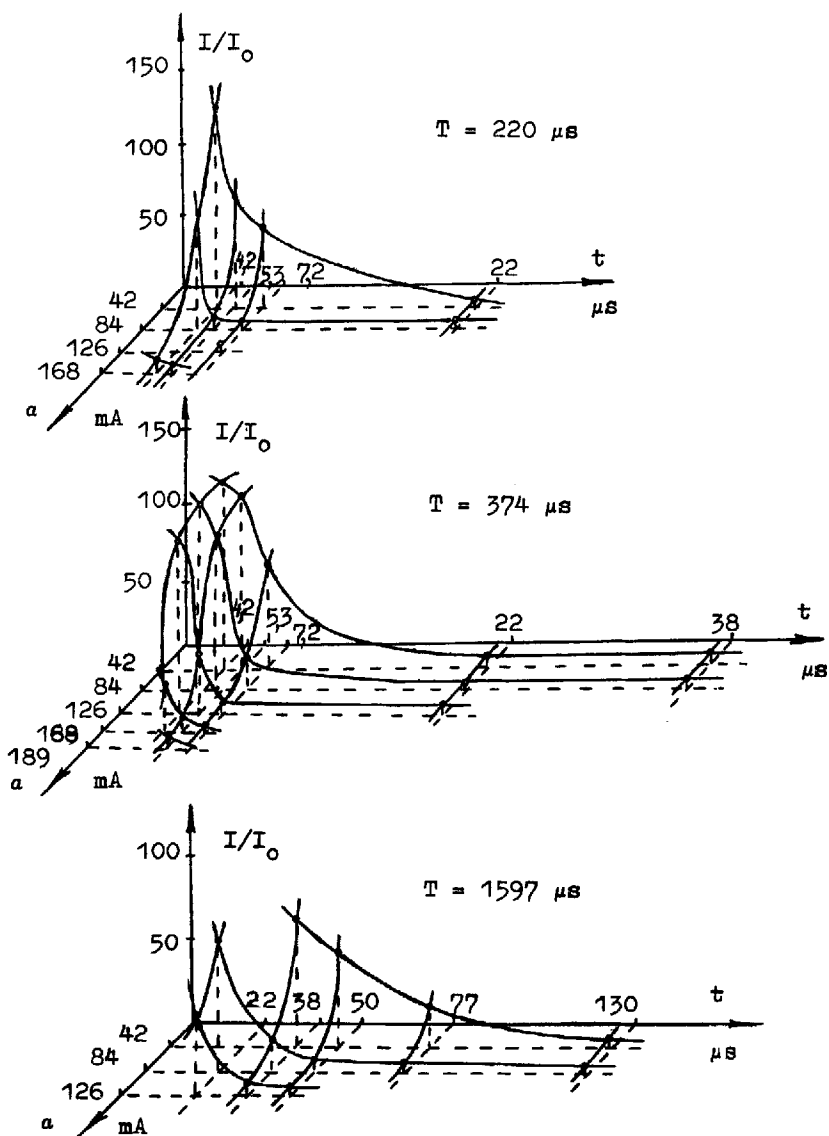


FIG.2

Dependence of Zn 330,2 nm spectral line intensity increase I/I_0 on the parameters of pulse supply.

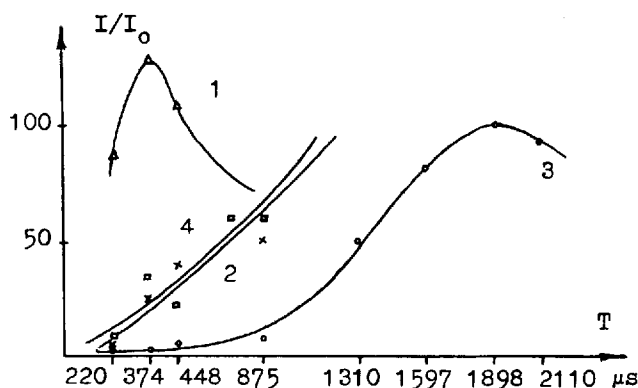


FIG.3

Dependence of Zn 330,2 nm spectral line intensity increase I/I_0 on T : 1- $t=4,2 \mu s$ at $a=84 \text{ mA}$; 2- $t=7,2 \mu s$ at $a=84 \text{ mA}$; 3- $t=22 \mu s$ at $a=84 \text{ mA}$; 4- $t=4,2 \mu s$ at $a=168 \text{ mA}$.

to higher values of t and a ; 5) I/I_0 decreases at $T \geq 1597 \mu s$. In tables 1 and 2 the widths of the Zn 330,2 nm spectral line emitted from dc. and pulse mode discharge correspondingly are summarized.

The Gaussian ($\Delta)_G$, $\Delta)_{G0}$) and the Dispersion ($\Delta)_D$, $\Delta)_{D0}$) components of the recorded widths ($\Delta)_0$, $\Delta)_{D0}$) are obtained according to Ballik method¹⁰. Corrections for Gaussian type broadening (due to the round diaphragm $\delta)_1$ and to the surface mirror defects $\delta)_2$) and of Dispersion type (as a result of the mirror reflection R - $\delta)_R$) were made by the equations given in¹¹:

$$\delta)_1 = \frac{r^2}{2F^2}, \quad \delta)_2 = \frac{4l}{\lambda}, \quad \delta)_R = \frac{1-R}{2n\lambda k},$$

where r is the radius of the round diaphragm, F is the focus distance of the collimating lens, l is the distance between the mirrors. Having in mind the corrections in our case ($\delta)_1=15 \text{ mK}$, $\delta)_2 \approx 24 \text{ mK}$, $\delta)_R=24 \text{ mK}$ at $R=90\%$, $l=0,7 \text{ sm}$, $r=0,05 \text{ sm}$, $F=50 \text{ sm}$, $\Delta l \approx \lambda/60$)

we obtain the proper Gaussian ($\delta\nu_G$, $\delta\nu_{Go}$) and the proper Dispersion ($\delta\nu_D$, $\delta\nu_{Do}$) components. Tables 1 and 2 indicate that:

- 1) $\delta\nu_D$ and $\delta\nu_{Do}$ do not depend on pulse parameters α , t and T ;
- 2) $\delta\nu_G$ and $\delta\nu_{Go}$ are 5+10 times higher than $\delta\nu_D$ and $\delta\nu_{Do}$. This coincides with the results obtained for the predominantly Gaussian character of the profile of the lines emitted from hollow cathode discharge¹¹. Besides it can be concluded that application of pulse supply gives rise mainly to thermal changes in the discharge. More clearly the dependence $\delta\nu_G = f(t, T)$ at $\alpha = 168$ mA and $\delta\nu_{Go} = f(i_o)$ /dashed line/ is shown on fig. 4.

Having in mind the results in tables and fig. 4, the following conclusions can be made: 1) $\delta\nu_G$ increases with α and t increasing; this increase is faster at $\alpha > 100$ mA; 2) $\delta\nu_G$ at $T = 220 \mu s$, $T = 374 \mu s$ and $T = 448 \mu s$ has approximately equal values, smaller or equal

TABLE 1

Dependence of Zn 330,2 nm widths $\Delta\nu_o$, $\Delta\nu_{Go}$, $\Delta\nu_{Do}$, $\delta\nu_{Go}$, $\delta\nu_{Do}$ [mK] on the discharge current i_o .

i_o [mA]	$\Delta\nu_o$	$\Delta\nu_{Go}$	$\Delta\nu_{Do}$	$\delta\nu_{Go}$	$\delta\nu_{Do}$
3,1	186,0	140,0	46,0	136,5	22,0
3,7	204,0	151,5	52,0	148,5	28,0
4,1	210,0	164,0	46,0	161,5	22,0
5,5	217,5	173,0	46,0	170,7	22,0
6,6	226,5	181,0	46,0	178,8	22,0
7,6	238,0	190,0	49,3	200,0	25,3
8,6	248,5	202,0	49,3	200,0	25,3
10,0	258,5	214,0	46,0	212,0	22,0

TABLE 2

Dependence of Zn 330,2 nm spectral line widths $\Delta\lambda$, $\Delta\lambda_G$, $\Delta\lambda_D$, $\delta\lambda_G$, $\delta\lambda_D$ [mK] on the pulse supply parameters (α , t, T).

T	μ s	α_{mA}	4, 2						5, 3						7, 2						22														
			$\Delta\lambda$	$\Delta\lambda_G$	$\Delta\lambda_D$	$\delta\lambda_G$	$\delta\lambda_D$	$\Delta\lambda$	$\Delta\lambda_G$	$\Delta\lambda_D$	$\delta\lambda_G$	$\delta\lambda_D$	$\Delta\lambda$	$\Delta\lambda_G$	$\Delta\lambda_D$	$\delta\lambda_G$	$\delta\lambda_D$	$\Delta\lambda$	$\Delta\lambda_G$	$\Delta\lambda_D$	$\delta\lambda_G$	$\delta\lambda_D$													
220		42	155	51	130	44	51	109	22	05	163	01	200	44	51	165	22	05	177	31	340	46	01	31	02	20	197	51	160	40	01	157	71	160	
		84	172	01	320	40	40	01	29	01	160	175	01	350	44	51	320	20	5	178	61	370	45	51	34	22	15	197	01	162	04	151	59	61	75
		126	175	01	340	44	51	131	02	05	180	51	430	41	51	40	31	75	183	21	450	42	81	42	71	188									
		168	190	51	480	46	01	45	72	20	195	51	550	41	51	52	31	75	220	01	1740	46	01	17	23	220									
374		42	150	31	107	04	45	103	02	05	161	21	150	46	01	114	220	174	31	320	44	51	129	02	05	195	51	152	04	45	14	97	20	5	
		84	160	01	190	41	151	115	71	75	173	81	310	44	51	280	20	5	177	41	350	44	51	132	62	05	169	31	160	04	00	157	41	160	
		126	168	81	270	41	151	23	81	75	179	51	390	42	81	360	188	182	21	430	42	81	40	31	188										
		168	189	51	470	44	51	144	32	05	194	01	1510	44	51	148	32	05	215	01	1700	46	01	16	77	220									
448		168	189	01	160	46	01	43	42	20	193	01	1500	46	01	147	32	20	218	01	1710	47	61	16	87	216									
875		168	190	01	1460	44	51	14	34	20	193	01	1490	44	51	14	60	20	5	209	01	1670	44	28	16	43	188	246	52	03	04	76	201	02	16
1897		168																									242	02	00	04	60	19	80	220	
2299		168																									251	02	09	04	60	207	02	20	

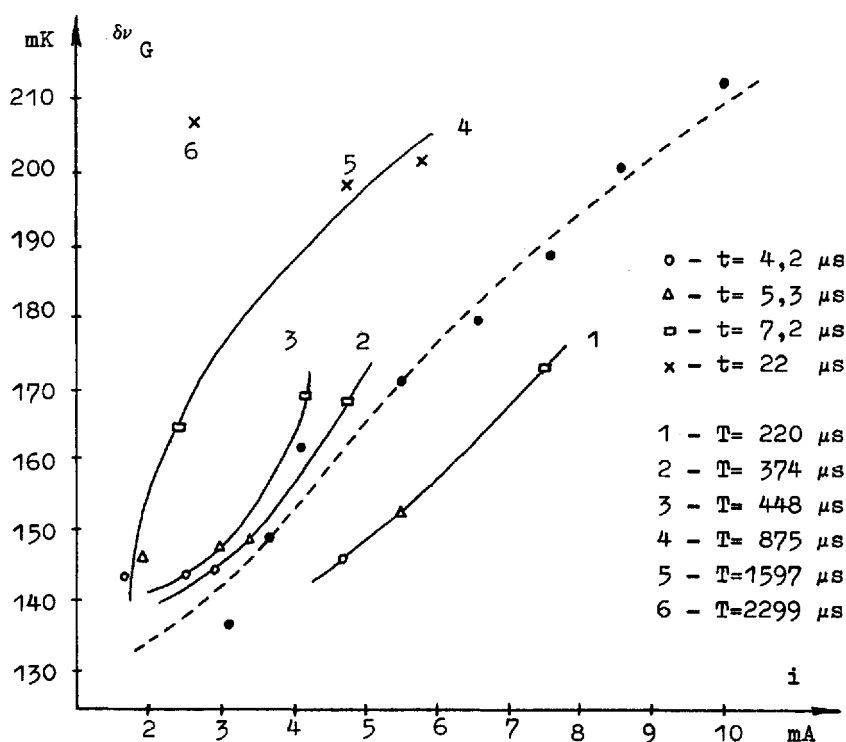


FIG.4

Dependences of proper Gaussian components $\delta\nu_G = f(t, T)$ at $\alpha = 168 \text{ mA}$ and $\delta\nu_{G0} = (i_0)$ /dashed line/ at the corresponding mean current values

to $\delta\nu_{G0}$; 3) $\delta\nu_G$ at $T \geq 875 \mu s$ increases and becomes bigger than $\delta\nu_{G0}$. While the first result can be explained by Doppler broadening and self-absorption because of mean current increasing in the discharge, such explanation for the second and the third results is not satisfactory. On the other hand our observations on the discharge at $T \geq 875 \mu s$ show that the shining diffused volume between the cathode and anode begins

to increase. Drawing a general conclusion from all these experimental results we can assume that at the given conditions the effect of the directed movement of the atoms from the cathode cavity outwards is manifested. Atomic transport phenomena from the cathode have been observed by other authors too^{12,13,14}. Detailed study of the investigated zinc line profile can lead to a more precise conclusion concerning the possibility of the assumed transport phenomena. The qualitative differences between our results and those for other elements obtained by the above cited authors regardless of some distinctions in the discharge tube constructions is predominantly due, in our opinion to the different properties of the investigated elements.

Generalizing the results obtained for the intensity and width of the Zn 330,2 nm spectral line emitted from hollow cathode discharge for our discharge tube construction in pulse mode, we can conclude that the optimum conditions for stable emission of zinc lines with increased intensity and at the same time with widths smaller or equal to that obtained at the d.c. mode are: $t=4,2+5,3 \mu s$, $a=40+80 \text{ mA}$ and $T=220+448 \mu s$.

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